Graphene Characterization using Time-Domain Terahertz Spectroscopy for Plasmonic Antenna Design

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ABSTRACT

Graphene's ability to support surface plasmon polaritons (SPPs) is of particular interest in the design of nanoscale plasmonic antennas. Since a dielectric-conductor interface is required to excite and sustain SPPs, a negative dielectric function becomes a defining property for graphene. We use terahertz time-domain spectroscopy (THz-TDS) to determine the complex dielectric function of graphene based on the extracted complex conductivity. These optical properties help us ascertain if a graphene sample is capable of supporting plasmons, and the appropriate dimensions to define a resonant cavity that would act as an antenna in the THz range.

CCS CONCEPTS

• **Hardware** → **Emerging tools and methodologies; Plasmonics;**

KEYWORDS

Graphene characterization, THz spectroscopy, plasmonic antennas

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1 INTRODUCTION

As wireless networking technologies continue to advance, our devices and systems are becoming more interconnected. At the same time, advances in materials science and nanotechnology are making possible increasingly smaller components with advanced functionality. Integrating these nanoscale components with an ever-evolving

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Figure 1: (Top) Ray diagram for THz light incident on the graphene-on-substrate system. (Bottom) Measured THz pulses corresponding to interface reflections.

network is expected to lead to an Internet of *Nano-Things* [1], which has the potential to revolutionize how we live and work.

Graphene boasts a long list of impressive physical properties [2]. Among these is the ability to support collective charge oscillations, known as *surface plasmons*, at a graphene-dielectric interface [3]. These surface plasmons can couple to electromagnetic (EM) waves, forming *surface plasmon polaritons* that strongly confine the free-space EM wave at the interface. Appropriate graphene dimensions would thus define a resonant cavity that should act as a plasmonic antenna in the THz range [5].

In this work, we use terahertz time-domain spectroscopy (THz-TDS) to extract the complex optical properties of CVD-grown graphene transferred onto an undoped silicon substrate. We do this by analyzing the light-matter interactions induced by picosecond pulses of THz light at the substrate-graphene interface as shown in Fig. 1. We measure a time-domain signal consisting of two pulses corresponding to the reflection at each interface. This is shown at the bottom of Fig. 1. We use a Hann window function centered at the peak of each pulse to separate the signal into two parts (red/blue in Fig. 1). We can then define a transfer function, $\tilde{H}$, using the echo pulse as the output and the main pulse as the input.

Figure 2: Extracted optical properties of graphene

2 THZ ANALYSIS AND CHARACTERIZATION

Our *self-referenced* measurement setup yields the reference substrate optical properties, $\tilde{n}_s = n_s - j\kappa_s$, in a single step through the transfer function equation defined below:

$$\tilde{H} = \frac{|\tilde{E}_2|}{|\tilde{E}_1|} \exp(-j[\phi_2 - \phi_1]) \quad (1)$$

Based on an approach similar to [4], we calculate the Fresnel reflection coefficient at the substrate–graphene–air interface as follows:

$$\tilde{R}_{sga} = \tilde{H}_g \frac{\tilde{R}_{as}}{\tilde{T}_{as} \tilde{T}_{sa}} \exp\left(j\frac{\omega}{c} 2\tilde{n}_s L_s \cos \theta_t\right) \quad (2)$$

We then extract the complex conductivity [8] using the equation:

$$\tilde{\sigma} = \frac{1}{Z_0} \left[\left(\frac{1 - \tilde{R}_{sga}}{1 + \tilde{R}_{sga}} \right) \frac{\tilde{n}_s}{\cos \theta_t} - \frac{1}{\cos \theta_i} \right] \quad (3)$$

Based on Maxwell's equations, the complex dielectric function is

$$\tilde{\epsilon} = \epsilon_0 + j\frac{\tilde{\sigma}}{\omega} = \epsilon_0 - \frac{\text{Im}[\tilde{\sigma}]}{\omega} + j\frac{\text{Re}[\tilde{\sigma}]}{\omega} \quad (4)$$

Since the real part of the dielectric function is related to the imaginary part of the complex conductivity, we fit our data, as shown in Fig. 2, using the Kubo formalism defined below, which is commonly used to model the intra-band surface conductivity of graphene [7].

$$\tilde{\sigma} = \frac{\tau}{1 - j\omega\tau} \left(\frac{2e^2}{\pi\hbar^2} k_B T \left[2 \cosh\left(\frac{E_F}{2k_B T}\right) \right] \right) \quad (5)$$

Here the Fermi energy of graphene is given by

$$E_F = \frac{\pi\hbar^2}{e^2} \frac{\sigma_{dc}}{\tau} \quad (6)$$

in which σ_{dc} is the dc conductivity and τ is the scattering time.

3 RESULTS

Based on the fitting to the Kubo formalism, we extracted the electrical parameters given in Table 1. We used the COMSOL Multiphysics platform to solve Eqs. (4) and (5) and simulate the electric field inside a graphene patch antenna. Figure 3a uses a calculated value of the scattering time [6], which shows the electric field of the surface plasmon being able to traverse the entire length of the patch. On

Table 1: Electrical parameters of graphene

Parameter	Symbol	Value
dc conductivity	σ_{dc}	1.3 mS
Scattering time	τ	31.5 fs
Fermi energy	E_F	0.35 eV
Carrier density	N	$1.4 \times 10^{13} \text{ cm}^{-2}$
Carrier mobility	μ	$576 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$

Figure 3: Electric field simulations for plasmon excitation in graphene patch antenna for different scattering times.

the other hand, Fig. 3b shows that if the extracted value of τ is used, then the electric field of the surface plasmon decays very quickly. A distinct advantage of THz–TDS over traditional electrical measurements is that, due to its non-contact nature, there is no risk of influencing or compromising the sample during measurement.

4 CONCLUSION

Here we report THz–TDS to be a robust optical and electrical characterization method for graphene, especially for applications involving plasmonic antennas. The clear difference seen in running simulations using predicted, or theoretical, values versus real-world values accentuates the importance of obtaining graphene parameters under real-world conditions to guide antenna design and better evaluate expected performance.

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